



Flow Behavior of Geldart A and Geldart C Particles in a Co-current Downflow Circulating Fluidized Bed Reactor

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ABSTRACT

The purpose of this research is to study the effect of Geldart A and C particles on the hydrodynamics behavior in a 9.3 m height, 0.1 m diameter co-current downflow circulating fluidized bed (downer reactor) using CFD simulation. Two-fluid model with kinetic theory of granular flow was adopted to predict flow behavior in the system. The simulation results show that hydrodynamics behavior in the downer strongly depends on the type of the particle. Geldart C particle exhibits a more uniform distribution along the lateral direction as compared with Geldart A particle. In addition, the effects of operating conditions were also studied. The uniformity of lateral direction of solids fraction increases with decreasing of solids circulation rate (G_s) or increasing of inlet superficial gas velocity (U_g). However, the radial distributions of gas and solids velocity are more uniform when U_g decreases especially for Geldart C particle.

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1. Introduction

A co-current downflow circulating fluidized bed (downer reactor) has been developed to overcome the drawback of a co-current upflow circulating fluidized bed or riser reactor. In the downer, both gas and solids particles are fed to the reactor at the top section. The suspension of gas and particle mixture co-currently downward flows along the gravitational direction. According to this feed setup, flow behavior in the downer is much more uniform as compared with that in the riser (Zhang *et al.*, 2001). Moreover, the particle residence time is lower and the particle residence time distribution (RTD) is narrower due to less back mixing (Khongprom *et al.*, 2012; Wei and Zhu, 1996).

Hydrodynamics behavior in the fluidized bed reactor strongly depends on the physical properties of particle (Khongprom, 2011; Khongprom *et al.*, 2012; Limtrakul *et al.*, 2008; Ye *et al.*, 2005). Small or fine particles are desired to apply in the chemical processes due to low mass and heat transfer resistances. Geldart A particles are normally applied due to its small particle size and can be easily fluidized at ambient conditions (Geldart, 1973). Geldart C particles are defined as an extremely fine particle. However, this particle type is rarely used in the conventional fluidized bed because these particles are very difficult to fluidize. According to feeding system and the assistant of gravitational force in the downer, Geldart C particles might be able to use in the downer reactor. Therefore, the objective of this research is to study the effect of Geldart A and C particles on the hydrodynamics behavior in the downer reactor.

2. Methodology

2.1 Reactor Geometry

The typical circulating fluidized bed (CFB) system is shown in Figure 1(a). This system mainly consists of downer section, gas-solids separator, riser section and gas and solids distributors. This work focuses on the hydrodynamics behavior in the 0.1 m inner diameter and 9.3 m height downer reactor (Figure 1(b)). Gas and solids particle are fed into the reactor at the top section and the exit locates at the bottom. To simplify the system, 2-D simulation was considered (see Figure 1(c)). Physical properties of particle used are shown in Table 1. Air at ambient condition was used as a fluidizing gas. The operating conditions used in this work are shown in Table 2.

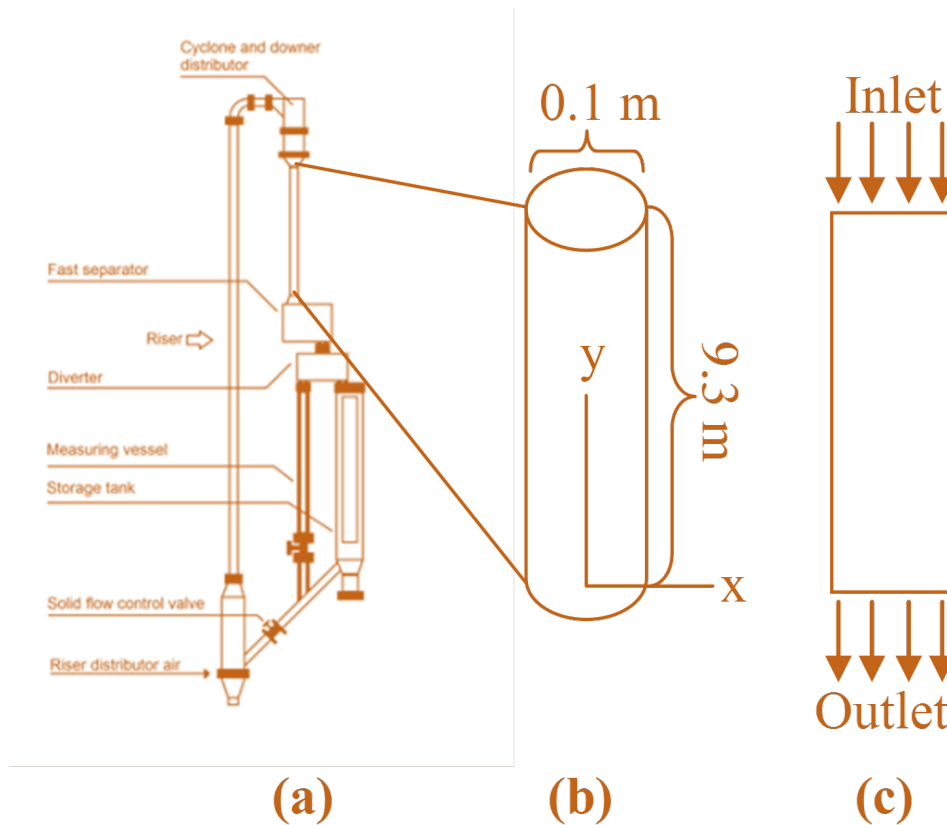


Figure 1: Typical CFB system (a); downer section (b); 2-D downer section (c).

Table 1: Physical properties of particle used in this study.

Particle type	D_p (μm)	ρ_p (kg/m^3)
Geldart A	80	1,500
Geldart C	20	1,500

Table 2: Operating conditions used in this study.

Case	U_g (m/s)	G_s ($\text{kg}/\text{m}^2\text{s}$)
Effect of G_s	3.5	150
		200
		250
Effect of U_g	2.5	101
	5.0	
	7.0	

2.2 Mathematical Modeling

Two-fluid model was adopted to predict flow behavior in the reactor. This model treats each phase as an interpenetrating continuum. Flow behavior of each phase was characterized by its own governing equations. The conservation equations are as follow,

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The continuity equation for phase q is expressed as:

$$\frac{\partial}{\partial t}(\alpha_q \rho_q) + \nabla \cdot (\alpha_q \rho_q \vec{u}_q) = 0 \quad (1)$$

where t stands for time (s), α_q the volume fraction, ρ_q the density (kg/m³), and \vec{u}_q the velocity vector (m/s).

Momentum conservation equation for gas phase is

$$\frac{\partial}{\partial t}(\alpha_g \rho_g \vec{u}_g) + \nabla \cdot (\alpha_g \rho_g \vec{u}_g \vec{u}_g) = -\alpha_g \nabla p + \alpha_g \rho_g \vec{g} + \nabla \cdot \overline{\overline{\tau}}_g + \nabla \cdot \overline{\overline{\tau}}_{tu,g} + \beta_{gp} (\vec{u}_p - \vec{u}_g) \quad (2)$$

where p is gas pressure (Pa), \vec{g} gravitational acceleration (m/s²), $\overline{\overline{\tau}}_g$ and $\overline{\overline{\tau}}_{tu,g}$ the viscous stress tensor and the Reynolds stress tensor, respectively (Pa), β_{gp} the interphase momentum transfer coefficient (kg/m³s).

Momentum conservation equation for solids phase (p) is

$$\frac{\partial}{\partial t}(\alpha_p \rho_p \vec{u}_p) + \nabla \cdot (\alpha_p \rho_p \vec{u}_p \vec{u}_p) = -\alpha_p \nabla p - \nabla p_p + \nabla \cdot \overline{\overline{\tau}}_p + \nabla \cdot \overline{\overline{\tau}}_{tu,p} + \alpha_p \rho_p \vec{g} + \beta_{pg} (\vec{u}_g - \vec{u}_p) \quad (3)$$

Interphase momentum transfer coefficient (Gidaspow's drag model) is defined as

$$\text{For } \alpha_g \leq 0.80: \quad \beta_{gp} = 150 \frac{(1-\alpha_g)^2 \mu_g}{\alpha_g d_p^2} + 1.75 \frac{(1-\alpha_g) \rho_g |\vec{u}_g - \vec{u}_p|}{d_p} \quad (4)$$

$$\text{For } \alpha_g > 0.80: \quad \beta_{gp} = \frac{3(1-\alpha_g)\alpha_g}{4} \frac{\rho_g |\vec{u}_g - \vec{u}_p|}{d_p} C_D \alpha_g^{-2.65} \quad (5)$$

$$\text{where } \text{Re} = \frac{\rho_g \alpha_g |\vec{u}_g - \vec{u}_p| d_p}{\mu_g} \quad (6)$$

where μ is gas viscosity (kg/m.s), d_p particle diameter (m), C_D drag coefficient (-).

Granular temperature conservation can be expressed in the form

$$\frac{3}{2} \left[\frac{\partial}{\partial t} (\rho_p \alpha_p \Theta_p) + \nabla \cdot (\rho_p \alpha_p \bar{u}_p \Theta_p) \right] = \left(-p_p \bar{I} + \bar{\tau}_p \right) : \nabla \bar{u}_p + \nabla \cdot (k_{\Theta_p} \nabla \Theta_p) - \gamma_{\Theta_p} + \phi_p \quad (7)$$

where Θ_p is granular temperature (m^2/s^2), \bar{I} unit tensor (-), k_{Θ_p} diffusion coefficient of granular temperature ($\text{kg}/\text{m}\cdot\text{s}$), γ_{Θ} collisional dissipation of solid fluctuating energy (kg/ms^3), ϕ_p energy exchange between phases (kg/ms^3).

The detail of this model has been discussed in previous works (Khongprom, 2011; Khongprom *et al.*, 2012).

k- ε turbulent model was used to explain the turbulent in the system. k equation for q phase is defined as

$$\begin{aligned} \frac{\partial}{\partial t} (\alpha_q \rho_q k_q) + \nabla \cdot (\alpha_q \rho_q \bar{u}_q k_q) = & \nabla \cdot \left(\alpha_q \frac{\mu_{tu,q}}{\sigma_k} \nabla k_q \right) + (\alpha_q G_{k,q} - \alpha_q \rho_q \varepsilon_q) + K_{pq} (C_{pq} k_p - C_{qp} k_q - \\ & K_{pq} (\bar{u}_p - \bar{u}_q) \cdot \frac{\mu_{tu,p}}{\alpha_p \sigma_p} \nabla \alpha_p + K_{pq} (\bar{u}_p - \bar{u}_q) \cdot \frac{\mu_{tu,q}}{\alpha_q \sigma_q} \nabla \alpha_q \end{aligned} \quad (8)$$

where k_q is the turbulent kinetic energy (m^2/s^2), G_k production of turbulent kinetic energy (kg/ms^3), ε_q turbulent dissipation rate (m^2/s^3), K_{pq} turbulent momentum transfer coefficient ($\text{kg}/\text{m}^3\text{s}$), σ_p turbulent Prandtl number.

Turbulent dissipation rate (ε) is expressed in the form

$$\begin{aligned} \frac{\partial}{\partial t} (\alpha_q \rho_q \varepsilon_q) + \nabla \cdot (\alpha_q \rho_q \bar{u}_q \varepsilon_q) = & \nabla \cdot \left(\alpha_q \frac{\mu_{tu,q}}{\sigma_\varepsilon} \nabla \varepsilon_q \right) + \frac{\varepsilon_q}{k_q} \left[C_{1\varepsilon} \alpha_q G_{k,q} - C_{2\varepsilon} \alpha_q \rho_q \varepsilon_q + \right. \\ & \left. C_{3\varepsilon} \left[K_{pq} (C_{pq} k_p - C_{qp} k_q) - K_{pq} (\bar{u}_p - \bar{u}_q) \cdot \frac{\mu_{tu,p}}{\alpha_p \rho_p} \nabla \alpha_p + K_{pq} (\bar{u}_p - \bar{u}_q) \cdot \frac{\mu_{tu,q}}{\alpha_q \sigma_q} \nabla \alpha_q \right] \right] \end{aligned} \quad (9)$$

where ε_q is turbulent dissipation rate (m^2/s^3), $C_{1\varepsilon}$, $C_{2\varepsilon}$, $C_{3\varepsilon}$ turbulent constants.

The kinetic theory of granular flow was used to predict the fluid properties of solids phase such as solids viscosity and solids pressure. The detail of this model was described somewhere else (Gidaspow, 1994).

2.3 Numerical Method

The governing equations were discretized using finite volume method. The first order upwind was applied for convection term. The SIMPLE scheme was used for solving the pressure and velocity in the system (Patankar, 1980). The 49 grid numbers in the radial direction and 140 grid numbers in the axial direction were employed. This grid system was obtained from grid independency study. A time step of 10^{-5} s was adopted. A convergence criterion of 10^{-4} for each scale was used.

3. Result and Discussion

3.1 Effect of Solids Circulation Rate (Gs)

Figure 2 shows the effect of Gs on the radial distribution of solids fraction at various heights. Both Geldart A and C particles exhibit core-annulus flow structure which uniform solids fraction in the center and high density peak near the wall region. This flow structure has been reported by several researchers (Cheng *et al.*, 1999; Khongprom, 2011; Khongprom *et al.*, 2012; Lehner and Wirth, 1999; Limtrakul *et al.*, 2008). However, Geldart C shows a much more uniform profile with smaller density peak near the wall. In addition, solids fraction increases with Gs.

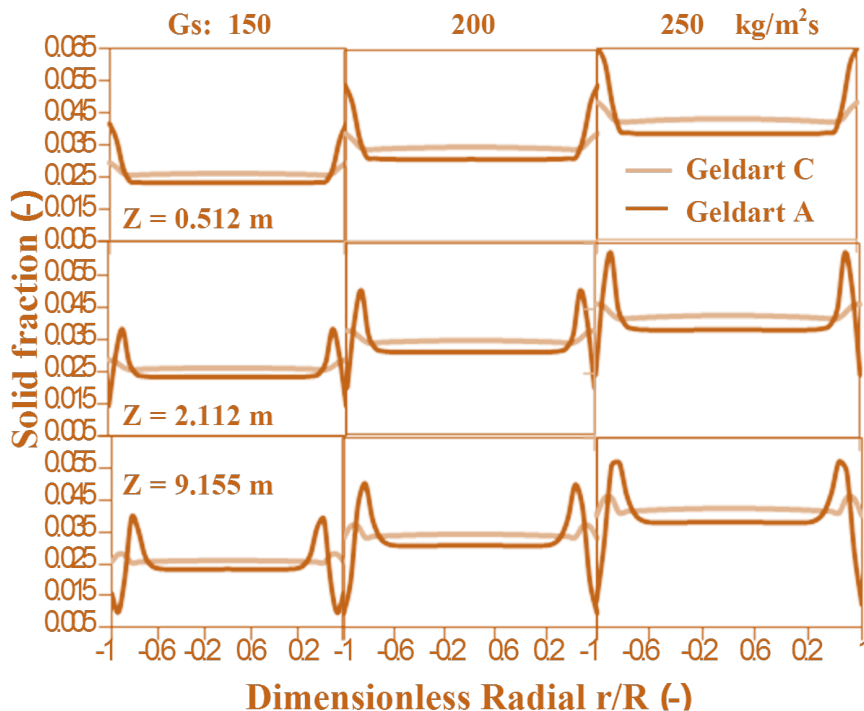


Figure 2: Effect of Gs on the radial distribution of time-averaged solids fraction.

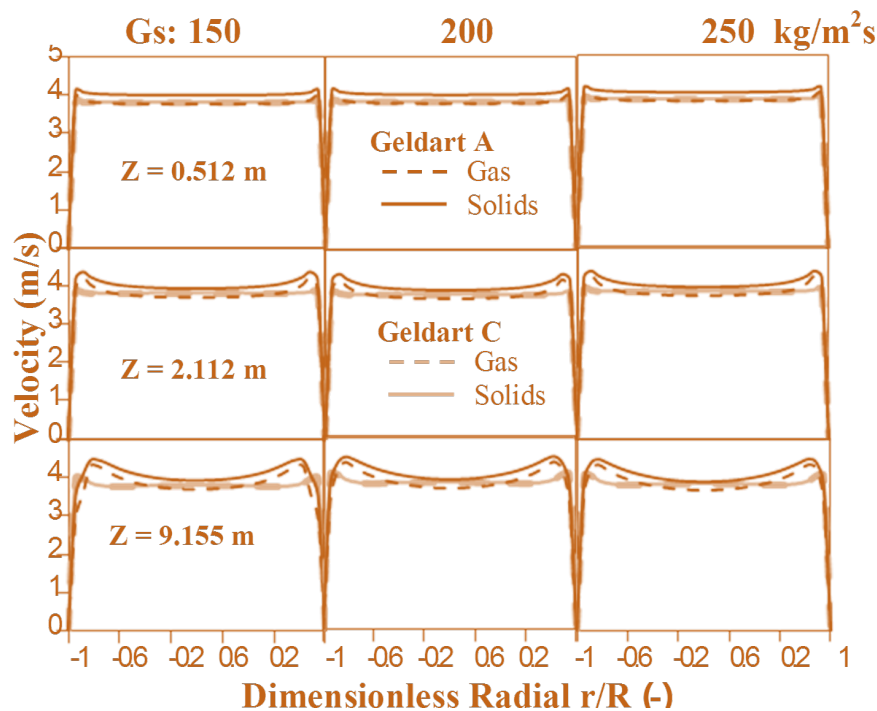


Figure 3: Effect of G_s on the radial distribution of time-averaged gas and solids velocities.

Figure 3 shows the effect of G_s on the radial distribution of gas and solids velocities. Near the inlet section ($Z = 0.512$ m), both gas and solids velocities are almost constant along the radial direction due to the uniform feed profile. Further down the column, Geldart C particle exhibits an almost uniform gas and solids velocities in the center region with a small velocity peak near the wall. Geldart A particle, both gas and solids velocities slightly increase from the center to the wall and the high velocity peak near the wall was form. This high velocity peak near the wall is the results of the particle cluster formation in this region (see Figure 2). These phenomena can also be observed by several experimental studies (Grassler and Wirth, 1999; Lehner and Wirth, 1999). In addition, the no slip velocity between gas and particles can be observed for Geldart C particle operation. This indicated that gas and Geldart C particle form a homogenous suspension mixture.

3.2 Effect of inlet Superficial Gas Velocity (U_g)

The effect of U_g on the lateral distribution of the solids fraction is shown in Figure 4. It was found that U_g exhibits less effect on the shape of the radial distribution profile of solids fraction. At $z = 0.512$ m, both Geldart A and C particles show a uniform profile in the center region with high solids fraction at the wall. Further down the column, the radial profiles

developed to form a core-annulus structure. However, solids fraction decreases with increasing of U_g . The density peak near the wall region decreases when U_g increases leading to more uniform in the radial direction especially for Geldart C particle.

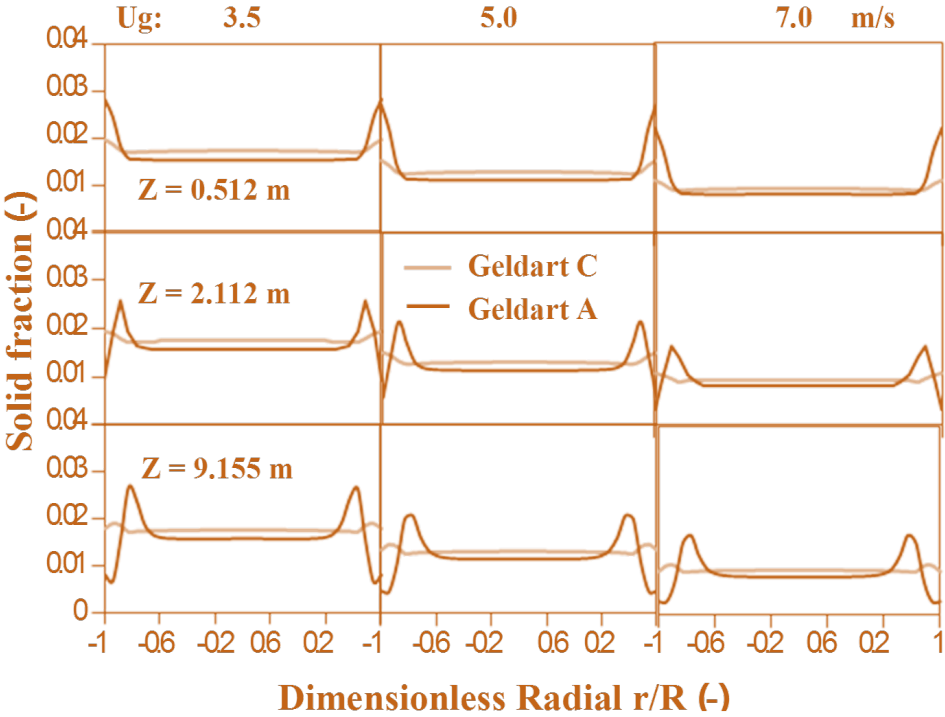


Figure 4: Effect of U_g on the radial distribution of time-averaged solids fraction.

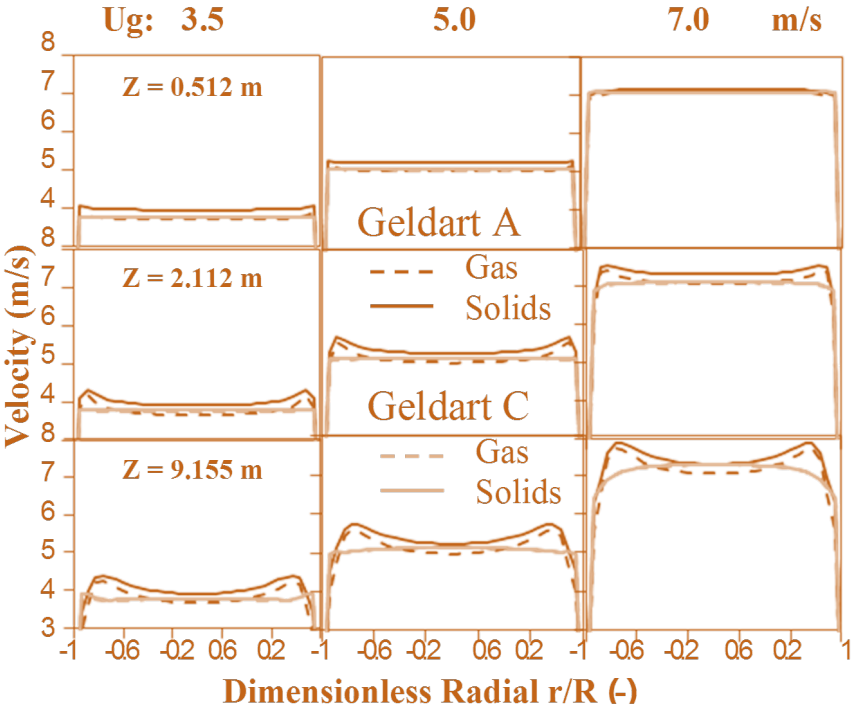


Figure 5: Effect of U_g on the radial distribution of time-averaged gas and solids velocities.

Figure 5 exhibits the effect of U_g on the radial profile of gas and solids velocities. As expected, gas and solids velocities increase with U_g . In case of Geldart C, gas and solids velocity profile in the fully developed region ($Z = 9.155$ m) can be classified into 2 types. At low U_g ($U_g = 3$ m/s), gas and solids velocities shows a uniform profile in the center with a small peak near the wall. At high U_g ($U_g = 7.0$ m/s), gas and solids velocities profiles shows a parabolic shape with consistency with gas velocity profile when operate with no solids particle feeding (did not show here). In this low solids fraction operating regime, gas phase governs the overall flow behavior in the system.

4. Conclusion

Two-fluid model with kinetic theory of granular flow was successfully developed to predict the hydrodynamics behavior in a downer reactor. The simulation results show that Geldart C particle exhibits a more uniform distribution along the lateral direction as compared with Geldart A particle. Geldart C particle exhibits a no-slip velocity between gas and solids particle phases. Moreover, high density peak near the wall region can be observed when operate with Geldart A particle. In addition, the effects of operating conditions were also studied. The uniformity of lateral distribution of solids fraction increases with decreasing of G_s or increasing of U_g . However, the radial distributions of gas and solids velocity are more uniform when U_g decreases especially for Geldart C particle.

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