


**PAPER ID: 10A11B**


## LONGITUDINAL FLOW OF A CONICAL DEFLECTOR WITH A STREAM OF VISCOUS INCOMPRESSIBLE LIQUID

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### ARTICLE INFO

#### Article history

Received 01 May 2019  
Received in revised form 09  
July 2019  
Accepted 25 July 2019  
Available online 31 July 2019

#### Keywords:

Jet flow; Conical nozzle;  
Water film thickness;  
Correction factor; Friction  
forces; Froude number;  
Trajectory; Jet departure  
distance.

### ABSTRACT

This article deals with a longitudinal jet axisymmetric flow around a conical deflector by an averaged turbulent flow of an incompressible fluid. An equation is given to determine the thickness of a water film that comes down from the deflector, and depends on the work of friction forces and kineticity of the incident flow. The distance of departure of the jet (radius of the deflector) and comparison with the experiment is given.

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## 1. INTRODUCTION

The longitudinal axisymmetric flow around a turbulent jet of a conical deflector is so complicated that at present. It is not possible to describe this process not only on the basis of the differential Navier – Stokes equations but also using the averaged Reynolds equations [1]. Wrapping conical deflectors is of great practical importance in the design of short-jet sprinkler heads. The flow of the fluid flowing around the conical deflector takes a fairly stable close to the conical shape of the film. It remains at some distance from the vanishing point of the surface of the deflector. After this, it loses stability and is broken up into separate drops. According to the experiments [2, 3], the disintegration of the jet begins from the peripheral part. Therefore, droplets forming in the central zone of the jet have the greatest range. The point of their fall is the point of maximum departure of the jet. The concept of “jet path” in this case refers to any axial section of an axisymmetric flow, with the assumption that the transverse dimensions of the compact part of the liquid film are small compared with the range and height of departure. In fact, it is believed that all the liquid moves along a single trajectory connecting the center of the living section of the film and the point of maximum departure, and is the envelope of the trajectories of all the droplets of the jet [2]. In the axisymmetric

case, the transverse dimensions of the compact part of a liquid film are uniquely determined by the thickness of the film that comes down from the deflector. Historically, the determination of the film thickness, as well as force interaction between the fluid and the deflector, is calculated without taking the friction forces into account [2-6]. In [7] it was established (Figure 1) that thickness of the film on the deflector depends on the full hydraulic drag coefficient  $\zeta$  (zeta), and Froude numbers of the incoming flow:

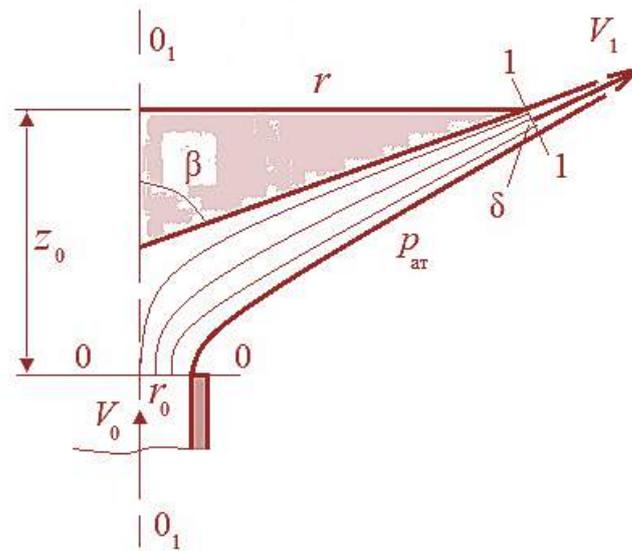
$$\delta = -\frac{r}{\cos\beta} + \sqrt{\frac{r^2}{\cos^2\beta} + k_0 \frac{r_0^2}{\cos\beta}} \quad (1),$$

where

$r, r_0, \beta$  – geometrical parameters of the deflector pair (nozzle - deflector),

$k_0 = \sqrt{\frac{\alpha_1 + \zeta}{\alpha_0 - \frac{2}{Fr_0}}}$  – correction factor that takes into account the work of friction forces and the

kineticity of the incident flow,  $\alpha_0, \alpha_1$  – Coriolis coefficients,  $Fr_0 = \frac{V_0^2}{gz_0}$  – Froude number.



**Figure 1:** Scheme of jet flow around a conical deflector.

The average velocities in the calculated living sections 0–0 and 1–1 are related by:

$$V_0 = k_0 V_1 \quad (2).$$

From Equation (1), it is inferred that the coefficient  $k_0$  can be determined empirically by the known geometrical parameters of the deflector pair and the film thickness.

Consider the trajectory and the departure length of the jet  $L$  (range nozzle). Neglecting the wind speed, let us imagine that the origin of coordinates is placed in the center of gravity of the living section 1 - 1 (see Figure 1). The distance from the center of gravity of the living section 1 - 1 to the

surface of the earth will be  $z_{\text{Hac}} + 0.5\delta\cos\alpha$ .

Insofar as  $z_{\text{Hac}} \gg 0.5\delta\cos\alpha$ , then we will later consider only  $z_{\text{Hac}}$  – deflector installation height.

With axisymmetric fluid motion in the deflector for any axial flow section in the absence of air resistance, the jet trajectory equation can be written as [8]:

$$z = xt g \alpha - \frac{g}{2V_1^2 \cos^2 \alpha} x^2.$$

Excluding speed  $V_1$  using equation (2), we get:

$$z = xt g \alpha - k_0^2 \frac{g}{2V_0^2 \cos^2 \alpha} x^2.$$

For cylindrical nozzle speed  $V_0$  can be represented by [8]:

$$V_0 = \mu \sqrt{2gH},$$

where  $\mu$  is the coefficient of speed (consumption) nozzle;  $H$  is head in front of the nozzle.

Then the jet trajectory equation takes as

$$z = xt g \alpha - \frac{k_0^2}{4\mu^2 H \cos^2 \alpha} x^2.$$

Energy loss to overcome air resistance when  $\alpha < \pi/4$  approximately equal to [8]:

$$h_f \approx \lambda \frac{l}{4R_1} \frac{V_1^2}{2g} = \frac{\lambda}{k_0^2} \frac{x}{4R_1 \cos \alpha} \frac{2g\mu^2 H}{2g} = \lambda \frac{\mu^2}{4k^2} \frac{Hx}{R_1 \cos \alpha},$$

where  $\lambda$  is drag coefficient;  $R_1$  is the hydraulic radius of section 1 - 1. Then, the final view of the jet's trajectory is written as:

$$z = xt g \alpha - \frac{k_0^2 x^2}{4\mu^2 H \left(1 - \frac{k_0^2 K}{4\mu^2 H \cos \alpha} x\right) \cos^2 \alpha} \quad (3),$$

where  $K = \lambda \frac{\mu^4 H}{k_0^4 R_1}$  is dimensionless parameter depending, in general, on  $r_0$ ,  $r_1$ ,  $\alpha$ ,  $\delta$ ,  $H$ , and resistance in the deflector and air.

From Equation (3) with  $z = -z_{\text{Hac}}$ , the radius  $L$  of the action nozzle can be determined from

$$-z_{\text{Hac}} = Lt g \alpha - \frac{k_0^2 L^2}{4\mu^2 H \left(1 - \frac{k_0^2 K}{4\mu^2 H \cos \alpha} L\right) \cos^2 \alpha} \quad (4).$$

The solution of the quadratic equation (4) has the following form [9]:

$$\left. \begin{aligned} L &= L_0 \cos \alpha, \\ L_0 &= z_{\text{Hac}} \frac{\sqrt{(aK - \sin \alpha)^2 + 4a(1 + K \sin \alpha)} - aK + \sin \alpha}{2a(1 + K \sin \alpha)}, \\ \alpha &= \frac{k_0^2 z_{\text{Hac}}}{4\mu^2 H}. \end{aligned} \right\} \quad (5)$$

Hydraulic radius  $R_1$  living section 1 - 1, which is a truncated cone with a generator  $\delta$  (see Figure 1) is approximately equal to  $\delta$ . Thus, the parameter  $K$  takes the form

$$K = \lambda \frac{\mu^4 H}{k_0^4 \delta} \quad (6).$$

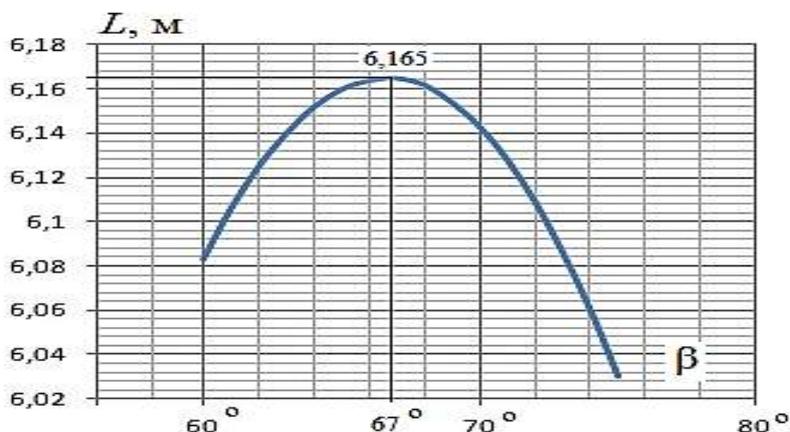
Friction coefficient  $\lambda$ , included in parameter (6), can be determined from the experimental data of A.I. Didebulidze [2].

According to calculations [7], the correction factor  $K_0$  is very weakly dependent on the Froude number  $Fr_0$ . Therefore, we further neglect the distance  $z_0$  (see Figure 1) and the Froude number.

The maximum value of  $L$  (the range of the nozzle), which depends on the angle  $\alpha$ , can be found numerically [10] using the algorithm (1), (5), (6), without solving the cumbersome extremum problem. Figure 2 shows a graph of dependence (4) according to [7, 11]:

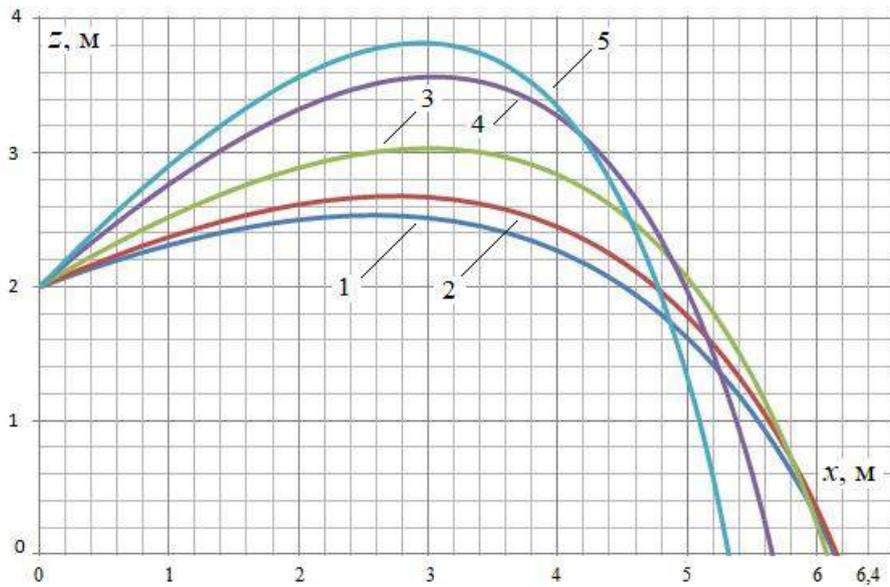
$$Q = 0.957 \text{ l/sec}; H = 10 \text{ M}; d_0 = 10 \text{ MM}; d = 50 \text{ MM}; \mu = 0.87; z_{\text{Hac}} = 2.0 \text{ M} \quad (7),$$

which shows that the maximum length of departure of the jet (range nozzle)  $L_{\text{max}} = 6.165 \text{ M}$  is achieved when  $\beta = 67^\circ$ , see Figure 2.



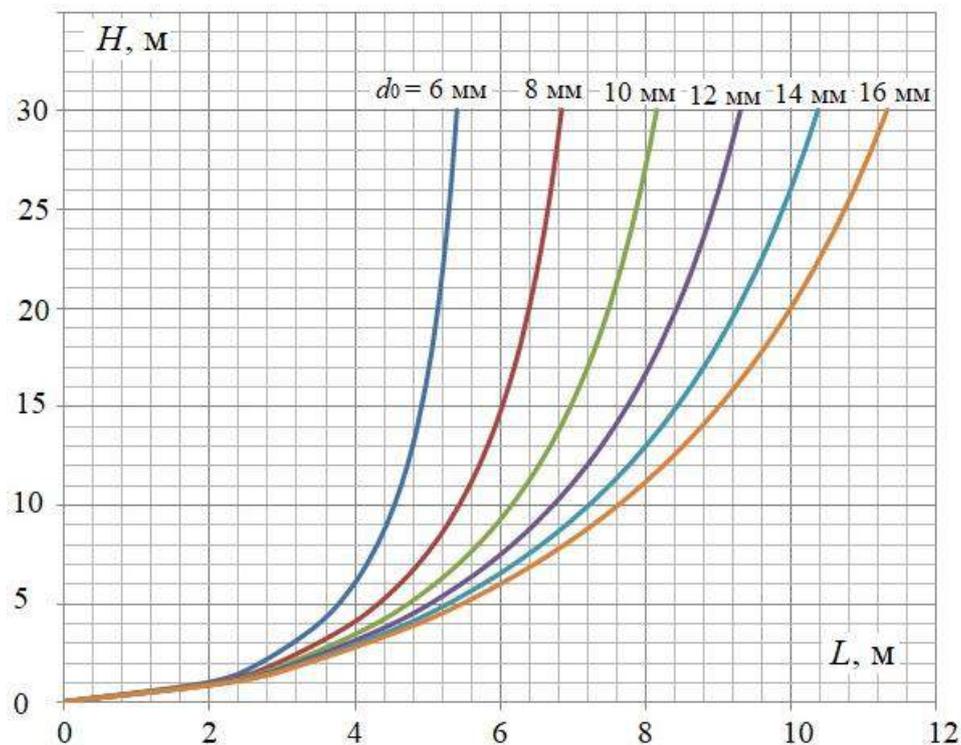
**Figure 2:** Determination of the range of action nozzle.

The calculation of the trajectories of the hydraulic jets according to equation (3) with the data (7) for different values of the deflection angle is presented in Figure 3.



**Figure 3:** The trajectory of the hydraulic jets at different values of the angle of taper: 1 –  $\beta = 70^\circ$ ; 2 –  $\beta = 67^\circ$ ; 3 –  $\beta = 60^\circ$ ; 4 –  $\beta = 50^\circ$ ; 5 –  $\beta = 45^\circ$ .

Using the algorithms (1), (5), and (6), it is possible to construct a theoretical dependence of the range of the nozzle on the head in front of the nozzle for different values of the nozzle diameter (Figure 4). Experimental confirmation of the curves in Figure 5 can be found in the work of N.F. Ryzhko [12].



**Figure 4:** Nozzle range at installation height 2 m.

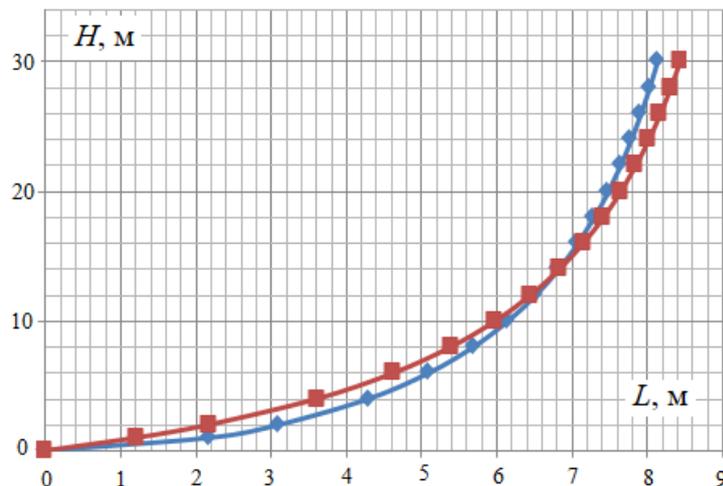
An empirical formula to determine the radius of action of a conical nozzle installed at a height of 2 m from the soil surface is also proposed there:

$$L = \frac{H}{0.728 + 0.942 \frac{H}{d_0}} \quad (8),$$

where  $H$  – head before nozzle in m,  $d_0$  – nozzle diameter in mm.

In Figure 5, there is a graph that compares the values of the nozzle range, calculated theoretically by the algorithm (1), (5), (6), as well as the empirical formula N.F. Ryzhko (8). Obviously, it is an enough good coincidence of the two dependencies.

Considering the specific features of conducting field experiments with operating sprinkler machines, the above studies allow us to conclude that some parts of the field tests can be replaced by computational experiments on appropriate models [10, 13].



**Figure 5:** Comparison of the calculated and experimental values of the nozzle range according to (7):  $\blacklozenge$  - theoretical calculation;  $\blacksquare$  - empirical formula (8)

Simulations of this kind [10] have been successfully used in various areas of the water sector, e.g. in water supply [14], and in hydraulic engineering [15-16].

## 2. CONCLUSION

This work focuses on a longitudinal jet axisymmetric flow around a conical deflector by an averaged turbulent flow of an incompressible fluid. The mathematical models are discussed to determine the thickness of a water film that comes down from the deflector, and depends on the work of friction forces and kineticity of the incident flow. The distance of departure of the jet (radius of the deflector) and comparison with the experiment is also discussed.

## 3. AVAILABILITY OF DATA AND MATERIAL

Data used or generated from this study can be requested to the corresponding author.

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